



A Coarse-Graining Procedure for Mapping Atomistic Models

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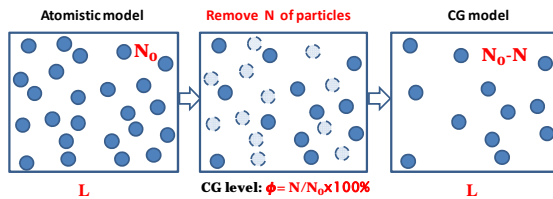
Motivation

Continuum mechanics: efficient but inadequate at nano scales.
Molecular dynamics: accurate but computationally expensive.
Goal: to develop effective hybrid computational methods for nano-scale flows.
Key issue: to determine whether an intermediate mesoscale description is required.

Model

Atomistic model: $N_0 = 1000$, ρ (number density) = 0.8, $r_{cut} = 2.5\sigma$ and Lennard-Jones (LJ) potential.

$$U^{LJ}(r) = 4\epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r} \right)^{12} - \left(\frac{\sigma_{ij}}{r} \right)^6 \right]$$



Coarse-grained (CG) model should preserve the thermodynamic and dynamic properties of the detailed atomistic system.

Method

Iterative Boltzmann inversion: construct an effective pairwise potential for the CG model that reproduces the radial distribution function, $g(r)$, of the detailed atomistic system.^[1]

$$U_0 = -k_B T \ln g_{target}(r)$$

$$U_{i+1}(r) = U_i(r) + \xi k_B T \ln \left[\frac{g_i(r)}{g_{target}(r)} \right]$$

Pressure matching: put pressure constraint on the effective potential through an optimization procedure.

➤ Objective function

$$f[U] = \frac{1}{2} \int_0^{r_{cutoff}} \left(\ln \frac{g_i(r, U)}{g_{target}(r)} \right)^2 dr + \frac{1}{2} \alpha \left(\frac{P_{CG} - 1}{P_{At}} \right)^2$$

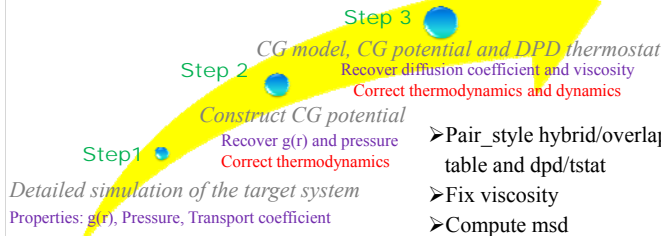
structure Penalty parameter Pressure (penalty function)

➤ Steepest decent algorithm

$$u_j^{k+1} = u_j^k - \frac{\partial f^k}{\partial u_j^k} \chi$$

Step size also match the pressure match g(r) Potential space

➤ Input CG potential table into LAMMPS



Dissipative Particle Dynamics thermostat (DPD)^[2]

➤ Galilean-invariant thermostat

$$\dot{p}_i = \vec{F}_i^C + \vec{F}_i^D + \vec{F}_i^R$$

➤ the dissipative force term mimics the friction between particles and dissipates energy.

$$\vec{F}_i^D = \sum_{j \neq i} \vec{F}_{ij}^D, \quad \vec{F}_{ij}^D = -\gamma w^D(r_{ij}) (\hat{r}_{ij} \cdot \hat{v}_{ij}) \hat{r}_{ij}$$

$$\vec{F}_i^R = \sum_{j \neq i} \vec{F}_{ij}^R, \quad \vec{F}_{ij}^R = -\sigma w^R(r_{ij}) \Theta_{ij} \hat{r}_{ij}$$

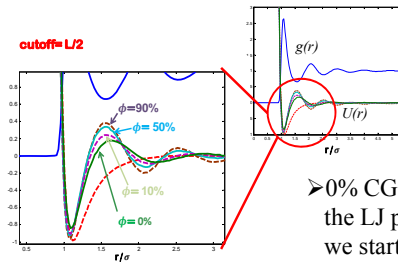
Viscosity: Reverse Nonequilibrium Molecular Dynamics (RNEMD)^[3]

$$j_z(p_x) = -\eta \frac{\partial v_x}{\partial z}, \quad j(p_x) = \frac{P_x}{2tA} \quad \text{➤ linear response theory}$$

Diffusion: Mean Square Displacement (MSD) $D = \frac{1}{6t} \sum_{a=1}^3 \langle [r_{ai}(t) - r_{ai}(0)]^2 \rangle$

Results and Discussion

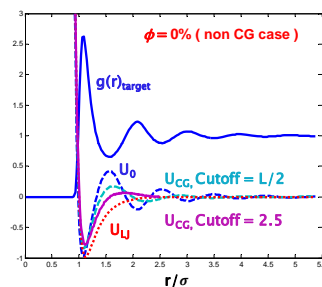
Effective potentials for various CG levels



➤ Potential of mean force is a good initial guess for high CG level case.

➤ 0% CG case doesn't reproduce the LJ potential ($r_{cut} = 2.5$) since we start with a longer-ranged interaction ($r_{cut} = L/2 \approx 5.4$).

Convergence test^[1]

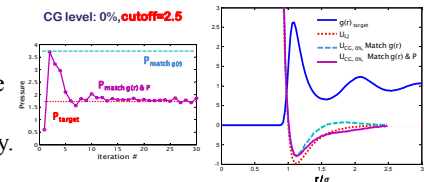


➤ The potential is closer to the LJ potential when r_{cut} is reduced to 2.5.

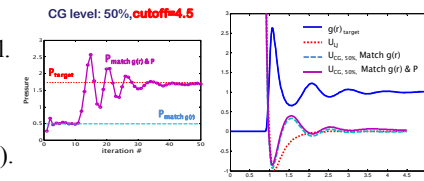
➤ the structure of a dense LJ ($\rho \approx 0.65$) system is mainly determined by the repulsive part of the potential.^[4]

After matching the pressure

➤ The attractive part of the potential is recovered. Now the potentials are more similar qualitatively.



➤ Pressure matching acts mainly on the long-range potential.

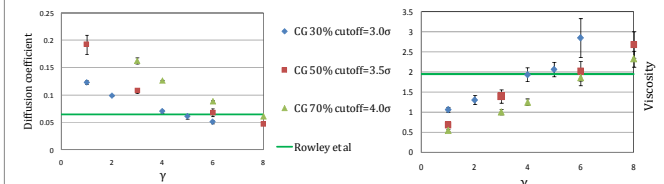


➤ Pressure is very sensitive to the potential, unlike g(r).

➤ Long-ranged interactions are important in high CG level case. The cutoff needs to be extended to recover g(r) and pressure.

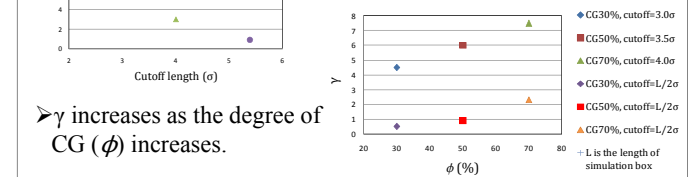
➤ The cutoff should be optimized based on the computational efficiency.

Match the dynamical properties



➤ A self-consistent value of γ for recovering both diffusion and viscosity is obtained for each CG level with the optimized cutoff length.

➤ The friction coefficient, γ , should increase as the length of interaction decreases.



➤ γ increases as the degree of CG (ϕ) increases.

References

[1] Reith, D. et al., *J. Comput. Chem.* 24, no. 13 (April 2003): 1624-1636.
[2] Groot, R. D. and Warren, P. B., *J. Chem. Phys.* 107, no. 11 (June 1997): 4423-4435.
[3] Muller-Plathe, F., *Phys. Rev. E* 59, no. 5 (May 1999): 4894-4898.
[4] Weeks, J. D. et al., *J. Chem. Phys.* 54, no. 12 (June 1971): 5237-5247.